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Density functional studies on photophysical properties and chemical reactivities of the triarylboranes: effect of the constraint of planarity

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Abstract The geometric and electronic structures, absorption spectra, transporting properties, chemical reactivity indices and electrostatic potentials of the planar threecoordinate organoboron compounds 1-2 and twisted reference compound Mes₃B, have been investigated by employing density functional theory (DFT) and conceptual DFT methods to shed light on the planarity effects on the photophysical properties and the chemical reactivity. The results show that the planar compounds 1-2 exhibit significantly lower HOMO level than Mes₃B, owing to the stronger electronic induction effect of boron centers. This feature conspicuously induces a blue shifted absorption for 1, although 1 seemingly possesses more extended conjugation framework than Mes₃B. Importantly, the reactivity strength of the boron atoms in 1-2 is much lower than that in Mes₃B, despite the fact that the tri-coordinate boron centers of 1-2 are completely naked. The interesting and abnormal phenomenon is caused by the strong p- π electronic interactions, that is, the empty *p*-orbital of boron center is partly filled by π -electron of the neighbor carbon atoms in 1-2, which are confirmed by the analysis of Laplacian of the electron density and natural bond orbitals. Furthermore, the negative electrostatic potentials of the boron centers in 1-2 also interpret that they are not the most preferred sites for incoming nucleophiles. Moreover, it is also found that the

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Department of Chemistry and Chemical Engineering, School of Chemical and Biological Engineering, University of Science and Technology, Beijing, China planar compounds 1-2 can act as promising electron transporting materials since the internal reorganization energies for electron are really small.

Keywords Density functional calculations · Laplacian · Photophysics · Reactivity · Triarylborane

Introduction

Organic conjugated oligomers and polymers have stirred great attention for several decades due to their wide applications in the field of organic light-emitting diodes (OLEDs), as well as fluorescent probes, nonlinear optics, organic field-effect transistors (OFETs) [1, 2]. The finetuning of the photophysical properties and emission characteristics has established them as an important class of new materials with intriguing properties. In these issues, the incorporation of the main-group elements into the organic conjugated materials provides an opportunity to modify the properties of the luminescent systems, since the interactions between the atomic orbitals of main-group elements and the π -conjugation system endow the organic materials with promising properties [3-8]. In this regard, one of the most appealing aspects is the incorporation of the electrondeficient boron element into organic conjugated structures, which has emerged recently as promising materials for OLEDs and nonlinear optics [9-14]. In addition to the four-coordinate boron complexes, the three-coordinate organoboron compounds have demonstrated their application potentials in the real word [15-22]. The availability of the empty p_{π} orbital on the boron center and the extended π system enable triarylboranes to show unusual optical and electronic properties. In resent years, many efforts have been devoted to designing and investigating the threecoordinate organoboron systems with novel properties

[23–25]. For instance, Yamaguchi and co-workers have recently shown that triarylboranes can be used as effective fluorescent sensors for Lewis donors such as fluoride [26, 27]. Shirota, Wang, Jäkle, and others have carried out extensive studies on the luminescent triarylboranes and have demonstrated their potential as effective emitters and electron-transporting materials [28–36].

In general, it is well known that the three-coordinate boranes are normally unstable toward hydrolysis by moisture. Therefore, the electron-deficient boron centers in the three-coordinate boranes need to be protected from external nucleophilic attack by using bulky substituents, such as phenyl, and 2,4,6-trimethylphenyl (mesityl) group. Kawashima and co-workers have synthesized several highefficiency ladder-type azaborines using mesityl substituent to protect tri-coordinate boron atoms [37]. Moreover, Wagner and co-workers recently synthesized 9,10-dianthryl-9,10-dihydro-9,10-diboraanthracene with 9-anthyl moieties protecting boron atoms, and proved their good stability toward air and moisture [38]. Recently, Yamaguchi et al. reported a group of planarized triarylboranes, which exhibits completely planar geometries and has no bulky substituents for protecting the naked boron centers [39, 40]. To our surprise, these compounds are very stable toward moisture regardless of the naked boron centers, which is in sharp contrast to the traditional triarylboranes. Piers et al. proposed that the enforced planarity of the triarylboranes may be an effective mean of stabilizing kinetically naked boron centers [41]. However, the origin of the stability of the planarized triarylboranes is still unclear for researchers. Therefore, it is essential to disclose the planarity effects on the photophysical properties and chemical reactivities of triarylboranes for an in-depth understanding of the structure-property relationships, thus designing stable functional materials.

Currently, quantum chemistry investigations have proved to be an important avenue in revealing the origin of chemical and physical properties and to assist in designing new materials with ideal properties [42-46]. Therefore, in this work, a systematically theoretical study was performed to reveal the planarity effects on the photophysical properties of triarylboranes, and to provide insight into the origin of the stability of the planarized triarylboranes. Specifically, the reported compounds 1 and 2 (Fig. 1) were chosen as models, and conventional twisted compound Mes₃B (Fig. 1) was also considered for reference and comparison. Their geometric and electronic structures, optical spectra, orbital components, ionization energies (IEs), electron affinities (EAs), internal reorganization energies (λ_{int}), Fukui functions, condensed and local softness, and electrostatic potentials were explored using density functional theory (DFT) and conceptual DFT calculations to provide insight into the structureproperty relationships and to understand the origin of the lowered reactivity of 1-2 toward nucleophiles.

Computational details

All calculations were performed using the Gaussian 09 package [48], unless otherwise stated. Comparative molecular structural calculations at the B3PW91/6-311+G(2df,2p) and B3PW91/6-31G(d,p) levels were carried out for 1. The largest difference between the bond lengths of the two levels is less than 0.005 Å, indicating that the 6-31G(d,p) basis set used in this work is accurate enough. Therefore, the equilibrium structures of the studied compounds were optimized based on the experimental crystal structures [39, 49] using the B3PW91/6-31G(d,p) level for the sake of time efficiency. Harmonic vibrational frequencies were calculated following the optimizations to ensure that real minima were obtained. All the calculations were carried out without any symmetry constraints. The absorption spectra were systematically calculated by TD-DFT method at B3PW91/6-311+ G(2d,p) level within the nonequilibrium polarizable continuum model (PCM) approach simulating the solvent effects (tetrahydrofuran, THF). In addition, the Fukui function [50, 51], local softness [52], dual descriptor [53], and Laplacian of the electron density were computed at B3PW91/6-31G(d,p) level through code Multiwfn 2.6 [54, 55] and their definition were summarized in Supporting information in detail. Generally, the adopted basis set affects the computed results of the Laplacian of the electron density. Luger et al. have demonstrated the influences of the basis sets on the calculated results of the Laplacian, *i.e.*, the more extended the basis is the more charge concentration is found in the bond. However, the computed values though different basis sets follow the same tendency within the studied compounds [56]. Therefore, the calculated results of Laplacian using B3PW91/6-31G(d,p) can provide reliable tendency for the studied compounds. The condensed descriptors, including condensed Fukui function, condensed softness, and condensed electrophilicity were calculated at B3PW91/6-31G(d,p) level using CHELPG scheme [57]. The electrostatic potentials (ESPs) for the studied compounds were computed at B3PW91/6-31G(d,p) level and were visualized using WFA Surface Analysis Suite [58]. The electronic configuration of boron centers were analyzed by natural bond orbitals (NBO) at B3PW91/6-31G(d,p) level.

Results and discussion

Geometric and electronic structures

The selected geometrical parameters optimized at the B3PW91/6-31G(d,p) level in gas phase together with the available X-ray crystal diffraction data are summarized in Table S1 (in Supporting information). The atomic labeling scheme is shown in Fig. 1. The relative errors between the



optimized bond lengths of Mes_3B , 1 and 2 and the corresponding experimental data are no more than 0.8 %, 0.7 %, and 0.7 %, respectively. Such discrepancies between the calculated and the experimental data are acceptable, indicating the B3PW91/6-31G(d,p) level can provide accurate prediction of the geometric structures. Notably, the bond length of B-C1 of 1 and 2 (1.523, 1.525 Å, respectively) is much shorter than that of Mes_3B (1.583 Å), indicating planarity strikingly impact the geometries of 1 and 2. Furthermore, the significantly shorter bond lengths for B-C may be indicative of the strong interaction of the atomic orbitals for boron and carbon atoms in 1-2 (*vide infra*).

To evaluate preliminarily the effect of constraint planarity on the electron distribution in frontier molecular orbitals (FMOs), Fig. 2 shows the plots of highest occupied molecular orbital (HOMO) and the lowest unoccupied molecular orbital (LUMO), energy levels and the energy gaps for Mes₃B, 1 and 2 at their optimized S₀ geometries by B3PW91/6-31G(d,p) level in vacuum. As shown in Fig. 2, the distribution of HOMOs and LUMOs for Mes₃B, 1 and 2 present similar features. The LUMOs mainly locate on the center boron atom and phenyl fragments, while HOMOs are somewhat localized on the phenyl fragments for the three compounds. Nevertheless, from the twisted compound Mes₃B to the planar compounds 1-2, their FMO energy levels are strikingly different. For 1, the LUMO level is similar to that of Mes₃B, which is consistent with the experimental result [39]. However, the HOMO level is significantly lowered as compared to that of Mes₃B, although the conjugation of 1 seems to be more extended than that of Mes₃B. For 2, the LUMO level is remarkably lowered, owing to the fact that 2 comprises two boron atoms which inherently hold the strong electron-withdrawing property. The HOMO level is slightly elevated compared

Fig. 2 The illustration of the orbital plots, the orbital energy levels, and the HOMO-LUMO energy gaps for compounds Mes_3B , 1 and 2 at their optimized S₀ geometries in vacuum at B3PW91/6-31G(d,p) level. The molecular orbital plots were obtained with an isosurface of 0.033 a.u. using ChemCraft [47]



with that of 1, since the conjugation plane is extended. However, the HOMO of 2 is still much lower than that of Mes_3B . The significantly lowered HOMO for 1-2 must have close relationship with the planarity effects, which can be apparently characterized by the perturbation to the electron cloud distribution of HOMO for 1. This interesting feature will be further revealed at the later section.

Absorption spectra and transporting properties

The absorption spectra calculated at PCM(THF)-B3PW91/6-311+G(2d,p) level were listed in Table 1. Moreover, to obtain more accurate results, a series of representative functionals including B3LYP [59, 60], B3PW91 [61-65], PBE0 [66-68], BMK [69], CAM-B3LYP [70], B971 [71], and M06 [72] were tested, taking 1 as an example. The calculated data and the experimental result collected in Table S2 (in Supporting information) show that B3PW91 provides the most accurate result, thus B3PW91 was adopted for the studied compounds. The calculated absorption wavelengths and the variation trends of the three compounds are in good agreement with experimental results, confirming that the adopted methods are reliable. The largest absorption lengths of the three compounds all originate from $S_0 \rightarrow S_1$ transition, mainly corresponding to the electron promotion from HOMO to LUMO. The largest absorption wavelength of 1 (313.3 nm) is blue shifted compared to that of Mes₃B (352.9 nm), demonstrating the perturbation of conjugation caused by planarity effects. 2 shows a strikingly red shifted absorption band than those of 1 and Mes₃B, which can be explained by its narrow band gap. Consequently, planarity effects exert pronounced influence on the optical properties of 1-2.

Ionization energies (IEs), electron affinities (EAs), and internal reorganization energies for hole (λ_{int}^+) and electron (λ_{int}^-) of **Mes₃B**, **1** and **2** have been calculated at B3PW91/6-31G(d,p) level to evaluate their transporting properties, and the results were listed in Table 2. From the computed results, the values of IEs for **Mes₃B**, **1** and **2** are relatively large, manifesting that it is rather difficult for hole injection due to the high barrier of injection [73]. Therefore, **1** and **2** can be used as potential hole-blocking materials in the OLEDs. In addition, the values of EAs for **1** and **2** are small, indicating

Table 2 The calculated vertical and adiabatic ionization energies (IEs) and electron affinities (EAs), internal reorganization energies for hole (λ_{int}^{+}) and electron (λ_{int}^{-}) of compounds **Mes₃B**, **1** and **2**. All values in eV

Compound	$I\!E_{\rm v}$	IEa	$EA_{\rm v}$	EAa	$\lambda_{\mathrm{int}}{}^+$	$\lambda_{\mathrm{int}}^{-}$
Mes ₃ B	7.288	7.158	0.556	0.447	0.260	0.212
1 2	7.644 7.347	7.583 7.296	0.376 1.054	0.417	0.124 0.106	0.083

the good ability of catch electron. Especially, it is important to note that the difference between the vertical and adiabatic IEs (EAs) for 1 and 2 is much smaller than that of Mes₃B. This fact implies that after gaining (or losing) electron, 1 and 2 undergo a smaller structural relaxation than Mes₃B, which is very important for the charge mobility. On the other hand, the planar configurations normally facilitate charge transport since they tend to prefer the tightly-stacked packing motifs and strong intermolecular electronic coupling [44, 74, 75]. Therefore, planar compounds 1 and 2 will facilitate the electron transport rather than Mes₃B. More importantly, from the computed results of internal reorganization energy, 1 and 2 both possess conspicuously smaller reorganization energies for electron (0.083, 0.116 eV for 1 and 2, respectively) than that of **Mes₃B** (0.212 eV). The values of λ_{int}^{-} for 1-2 are even lower than those of perylene bisimide (PDI) derivatives (ca. 0.3 eV, obtained at B3LYP/6-31G(d,p) level) [76], which are typical electron transporting materials. This conceivably supports that 1 and 2 could act as better electron-transport materials. However, three-coordinate boranes are normally unstable due to the reactivity of the electron-deficient boron atoms. Therefore, the chemical reactivities of the 1-2 and the reference compound Mes₃B were characterized in detail in the next section.

Reaction activity

In recent years, within the context of DFT, many useful and important reactivity indices, such as Fukui function, electronegativity, hardness, and softness, have provided reliable interpretations and predictions of chemical reactivity sites and the mechanism of the chemical reactions [77–86]. To

Table 1 Calculated absorption (λ_{abs} , nm) wavelengths, excitation energies (E_x , eV), oscillator strengths f, and dominant excitation character of **Mes₃B**, **1** and **2** together with experimental results ($\lambda_{exp.}$, nm).

Calculations were performed at PCM(THF)-TD-B3PW91/6-311+G(2d,p)//B3PW91/6-31G(d,p) level

Compound	Transition	λ_{abs}	$E_{\rm x}$	f	Composition ^a	λ_{\exp}^{b}
Mes ₃ B	$S_0 \rightarrow S_1$	352.9	3.513	0.136	H→L (98 %)	332
1	$S_0 \rightarrow S_1$	313.3	3.957	0.108	H→L (95 %)	320
2	$S_0 \rightarrow S_1$	370.4	3.347	0.212	H→L (97 %)	377

^a H denotes HOMO and L denotes LUMO ^b Measured in THF, ref. [39]

find the most preferred site for nucleophilic attack, the dual descriptor Δf , Fukui function and local softness toward nucleophiles (f^+ and s^+) were calculated at B3PW91/6-31G(d,p) level using code Multiwfn 2.6 [54, 55] and the results were illustrated in Fig. 3. All the plots of Δf , f^+ , and s^+ show that the maximum site is on the boron atoms for the three compounds, demonstrating the reactivity of boron atoms for nucleophilic attack. However, the contribution from the boron center to the Fukui function and local softness clearly decreases from Mes₃B to 1-2, which is indicative of the lowered reactivity toward nucleophiles for the latter two compounds. This is in good agreement with the experimental results. Furthermore, it is observed from the plots of Δf that **Mes₃B** shows evidently negative Δf (favored for electrophilic attack) regions at C1 species, while those in 1-2 are remarkably small, suggesting Mes₃B also shows relatively stronger reactivity toward electrophiles.

To more clearly compare the chemical reactivity strength of the boron centers, a series of condensed descriptors for reactivity, including condensed Fukui function, condensed softness, and condensed electrophilicity has been computed using CHELPG scheme [57]. Generally, the condensed reactivity descriptors are sensitive to the adopted method and population partitioning scheme. From the previous works, the choice of DFT methods based on exchange correlation functional is crucial for the calculation of the condensed descriptors [78]. In 2002, Subramanian et al. had investigated the effects of the basis set and population scheme on the results of the condensed Fukui functions [87]. Their results showed that CHELPG is a reliable scheme and could provide precise reactive site with less dependency on the basis sets. Therefore, all the condensed descriptors were calculated at B3PW91/6-31G(d,p) level for the sake of computational cost. The condensed softness, which combines local (atomic) and global information for a molecule, is strikingly important for the interpretation of reactions for the same sites in different molecules [88]. It has been found that for frontier orbital controlled softsoft interactions, the increasing condensed softness of the contact atom will result in an increase of the chemical reactivity [81]. From Table 3, evidently, the reactivity of the boron atom for Mes₃B is much higher than those of 1 and 2, since the values of the condensed softness are much larger for boron center in Mes₃B. The reactivity of the boron atoms for nucleophilic attack decreases in the order of $Mes_3B > 1 > 2$. This feature can also be reflected by the atomic charge values (q) for the boron atom. The atomic charge of Mes₃B is more positive than 1-2, demonstrating the stronger electrostatic attraction to the incoming nucleophiles. Moreover, Mes₃B presents a condensed electrophilicity value of $\omega^+(B) = 0.593$ at boron center. However, the constraint planarity decreases the condensed electrophilicity power of



Fig. 3 The maps of dual descriptor (Δf) , Fukui function (f), and local softness (s) toward nucleophilic attack for Mes₃B, 1 and 2. The isosurfaces surface of Δf , f^+ , and s^+ is 0.0032, 0.0032, and 0.0003 a.u., respectively. For the plots of Δf , purple is positive zone and blue is negative zone

			-						
Compound	η	μ	ω	S	$q(\mathbf{B})$	$f^{+}(B)$	<i>s</i> ⁺ (B)	$\omega^+(\mathrm{B})$	
Mes ₃ B	6.841	-3.868	1.093	0.146	0.710	0.542	0.079	0.593	
1	7.268	-4.010	1.106	0.138	0.344	0.457	0.063	0.506	
2 ^a	6.293	-4.201	1.402	0.159	0.307	0.243	0.039	0.341	

Table 3 The computed results of chemical hardness (η , eV), chemical potential (μ , eV), electrophilicity index (ω , eV) and global softness (S, eV⁻¹) for **Mes₃B**, **1** and **2**. Condensed Fukui functions (f^+ , a.u.),

softness (s^+ , a.u.eV⁻¹), and electrophilicity (ω^+ , a.u. eV) of the contact boron atom using CHELPG method are also listed. All the values were computed at B3PW91/6-31G(d,p) level

^a For 2, the related condensed values for B2 are same to B1

boron center from 0.593 in **Mes₃B** to 0.506 in **1**. Note that the condensed local electrophilicity for boron atom is mainly derived by change in the condensed Fukui function. Consequently, the boron centers of **1-2** show relatively lower reactivity toward nucleophiles compared to that of **Mes₃B**, which agree well with the experimental findings.

It is well-known that the three-coordinate boranes normally need bulky substituents to protect the electrondeficient boron atom from nucleophilic attack. Surprisingly, the boron centers in 1-2 are more stable toward nucleophilic attack than that in Mes₃B according to the calculated results, although the empty p_{π} orbital of the boron center is completely naked. To reveal the origin of this interesting and abnormal phenomenon, we will next look at the Laplacian of the electron density and electrostatic potential.

Laplacian of the electron density and electrostatic potential

To unveil the origin of the lowered reactivity of the boron centers in 1-2, the components of their LUMOs were first analyzed and the result was plotted in Fig. 4. As discussed



Fig. 4 The lowest unoccupied molecular orbital (LUMO) components (including three parts: boron, phenyl, and others, in percent) for compounds Mes_3B , 1 and 2 at optimized S_0 geometries in vacuum at B3PW91/6-31G(d,p) level

above, the boron center makes negligible contribution to the HOMO, while its contribution to LUMO is strikingly different for Mes_3B , 1 and 2 (see Fig. 2). Figure 4 shows that the contribution of the boron is greatly reduced from Mes_3B to 1. Turning from 1 to 2, the contribution of the boron atoms to the LUMO is slightly higher, owing to the fact that 2 has two boron atoms. Note that the contribution of boron for 2 is still smaller than that for Mes_3B , again proving that the planar configurations facilitate the extension the π -conjugation electron and stabilize the empty *p* orbital of the boron center.

Laplacian of the electron density $\nabla^2 \rho(r)$ is a useful tool for characterizing bonding interactions [89–94]. The investigation of the Laplacian of electron density could also provide insight into the electronic properties of atoms and molecules, and allow for useful predictions of chemical reactivity [95–99]. In general, regions where the Laplacian is negative reflect the charge is locally concentrated, which are characteristic of covalent bonding interactions. Conversely, regions where the Laplacian is positive indicate the charge is depleted, which are typical of closed-shell interactions. For the sites with negative Laplacian, they tend to deliver electronic charge and reactive toward electrophilic reactants, while those with positive Laplacian, they are apt to accommodate extra electronic charge and reactive toward nucleophilic reactants [93].

The Laplacian of the electron density shown in Fig. 5 has been computed at B3PW91/6-31G(d,p) level using code Multiwfn 2.6 [54, 55]. In addition, Wiberg bond index (WBI), which is a useful description of delocalization [100-102], was also computed to gain further insights into the bonding situation and the results were collected in Table S3 (Supporting information). It is readily seen from the plots of Fig. 5 and Fig. S1 that the carbon atoms of the phenyl ring are sp^2 hybridized and the electron-rich regions locate around carbon atoms for Mes₃B, 1, and 2. The electrondepleted region (electron-poor zone) is found in center boron atom, confirming the electron-deficient characteristic of boron center. It is evident that the boron center is stabilized by the charge-concentrated neighbor carbon atoms. Interestingly, the extension of the charge concentration region in B-C1 bond for 1 and 2 is much greater than that for Mes₃B,



Fig. 5 Contour maps of the Laplacian of the electron density $(\nabla^2 \rho)$ for Mes₃B (*top*) and 1 (*bottom*), in the plane defined by atoms B, C1 and C2 (The Laplacian of 2 resembles that of 1, hence the related plots was presented in supporting information). The dash lines denote negative Laplacian, while the solid lines denote positive Laplacian

manifesting that the interactions between boron center and neighbor carbon atoms are much stronger for 1 and 2. The relatively stronger interatomic orbital interactions in 1-2 also can be indicated by the larger values of WBI for B-C1 (0.949, 0.951, 0.888 for 1, 2, Mes₃B, respectively). More importantly, the delocalization of π electron in C1 atoms into the vacant *p* orbital in boron center can be further proved by the calculation results of natural bond orbitals (NBO) analysis. The electronic configuration of boron atom is 2s(0.62)2p(1.47) for Mes₃B, while those for 1-2 are 2s(0.57)2p(1.68), 2s(0.58)2p(1.65) (B1 in 2), 2s(0.58)2p(1.65) (B2 in 2), respectively. The net electrons for boron atoms in 1-2 are clearly larger than that in Mes₃B, suggesting that the empty p_{π} orbital has been partly filled for 1-2. Therefore, the stabilization effects provided by chargeconcentrated C1 species are much greater in 1-2 than those in **Mes₃B**. This feature is strikingly important for the lowered reactivity toward nucleophiles for 1-2, whose boron centers is naked without any protection of bulky substituents and shows higher stability toward nucleophilic attack.

For carbon atoms C7 and C8 in 1 and 2, the extension of charge concentrated zone for C7-C3 and C8-C2 is smaller than the bonds of the phenyl ring (such as C2-C4, C4-C6, etc.). The related bonds associated to C7 and C8 do not participate in the conjugation system, since they have single bond character as manifested by WBI values which are close to 1.0. Importantly, for the three compounds the extension of charge concentrated region in the bonds of C1-C2 and C1-C3 is different from those of C2-C4, C3-C5, C4-C6, and C5-C6, where the former are relatively smaller than the later. Moreover, the computed nucleus independent chemical shift (NICS) values at the geometric center phenyl rings calculated at B3LYP/6-31+G(d,p) level for 1 and 2 (-6.78 and -6.57) are much higher than that of benzene (-8.04) [103, 104]. These characteristics manifest boron atoms which strikingly disturb the conjugation system. The chargedepleted boron atoms in 1 and 2 clearly have no contribution to the conjugation system. In converse, the empty p_{π} orbital of boron center needs to be stabilized by the conjugation system. Coincidentally, the planar configuration of 1 and 2 is beneficial for the boron center extract electron from



Fig. 6 The electrostatic potential (in a.u.) of compounds Mes_3B , Ph_3B , 1 and 2 calculated at B3PW91/6-31G(d,p) level. The electrostatic potential values are mapped on the surface of electron density (0.001 a.u.) and created with WFA Surface Analysis Suite [58]. Color range: blue <-0.017, green -0.017 - 0, yellow 0 - 0.010, red >0.010

neighbor carbon atoms. This also can be characterized by the relatively smaller value of the EA_v for 1 (0.376 eV for 1 and 0.556 eV for Mes₃B), which suggests the electrondeficient boron atom has been stabilized. Therefore, the HOMO level of 1 is lowered compared to that of Mes₃B, although they possess similar conjugation framework (vide supra). In addition, to clarify the contribution of the conjugation and inductive effects to the lowered HOMO of 1 and 2, the simplest arylborane compound, $C_6H_5BH_2$, was calculated. The planar and twisted (the BH₂ subunit is perpendicular to benzene plane) C₆H₅BH₂ model and the calculation details were presented in the Supporting information (Fig. S2). The HOMO level is -7.14 and -6.64 eV for the planar and twisted model, respectively, which exhibit similar tendency as 1 and Mes₃B. For the planar model, the electronic configuration of boron atom is $2s(0.85)2p_x(0.71)$ $2p_{\nu}(0.90)$ $2p_{z}(0.14)$, while the twisted model is 2s $(0.87)2p_x(0.72) 2p_v(0.02) 2p_z(0.88)$. This feature indicates that the conjugation effects of the boron center are more important for the lowered HOMO.

Additionally, within the framework of hard and soft acids and bases of Pearson [105, 106], the electrostatic potential (ESP), which is frequently used to predict the reactive sites for hard-hard interactions and the packing motifs of a crystal [107–110], has been computed (at B3PW91/6-31G(d,p) level in vacuum) for Mes₃B, 1 and 2. The most positive (negative) regions of ESP are generally preferred for nucleophilic (electrophilic) attack. From the plots of Fig. 6, the hydrogen atoms present the most positive ESP regions, while the conjugated systems (phenyl rings) show the most negative regions for the three compounds. Interestingly, the boron centers of the planar compounds 1-2 take negative electrostatic potential values, which is less negative than those of conjugated systems. This proves that the boron centers are not the most preferred sites for nucleophilic attack, in line with the above discussions of the reactivity descriptors. It is well-known that the boron center of Mes₃B is well protected by methyls in mesityl groups. The positive ESP regions near the boron center in Mes₃B are probably attributed to the hydrogen atoms in methyl groups. Therefore, the ESP of the triphenylboron (Ph₃B in Fig. 6) was also computed to distinctly compare the ESP behaviors of the boron centers. From Fig. 6, the boron center of the Ph₃B clearly presents a positive ESP value, demonstrating the high reactivity toward nucleophiles. This further confirms that the planarity effects strikingly reduce the strength of the chemical reactivity for boron atoms toward nucleophiles.

Conclusions

The aim of the present investigation was to assess the planarity effects on the photophysical properties and the chemical reactivities of three-coordinate organoboron compounds 1-2 as compared to the twisted reference compound Mes₃B. To this end, within the framework of DFT and conceptual DFT the geometric and electronic structures, absorption spectra, ionization energies and electron affinities, chemical reactivity indices and electrostatic potentials have been calculated. The results can well reproduce the experimental measurements. With the planarity configurations, 1 exhibits significantly lower HOMO level than Mes₃B, owing to the stronger electronic induction effect of boron center. The characteristic directly induces a blue shifted absorption for 1 compared with Mes₃B, although 1 seemingly possesses more extended conjugation framework than Mes₃B. 1-2 can act as promising electron transporting materials since their reorganization energies for electron are really small. Interestingly, the completely naked boron atoms of 1-2 are more stable toward nucleophiles than the well protected one in Mes₃B according to the calculated chemical reactivity indices. This interesting and abnormal phenomenon can be explained by the stronger electronic interactions between boron center and neighbor carbon atoms in 1-2, which are confirmed by the analysis of Laplacian of the electron density and NBO. Furthermore, the negative electrostatic potentials of the boron centers in 1-2 also rationalized the lowered chemical reactivity.

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